

Notes for ECE-606: Spring 2013

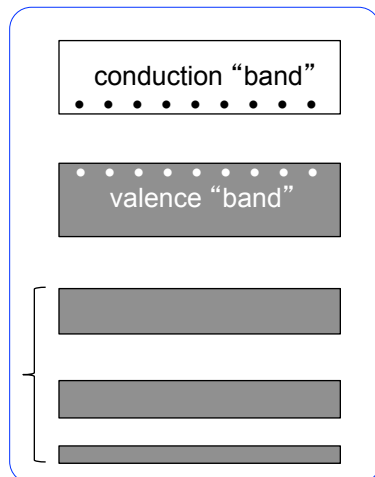
DOS(E) & Carrier Densities

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Si energy levels / bands

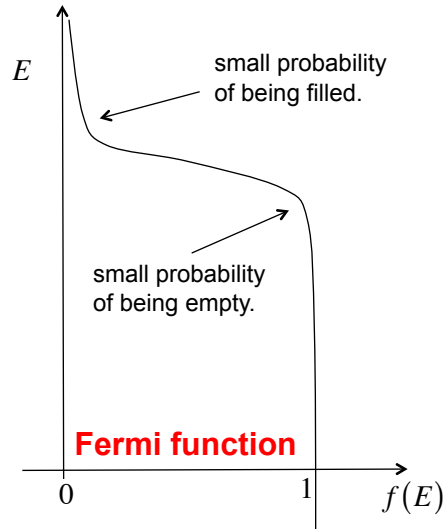
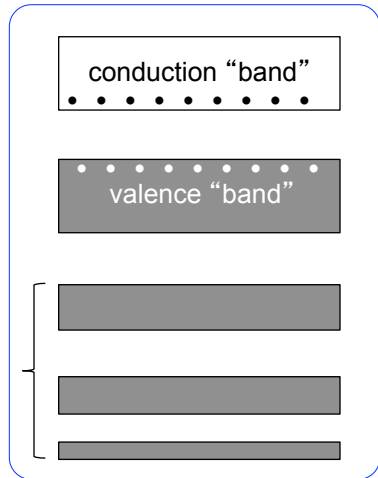


***What is the probability, f ,
that the states in any of
these bands are filled?***

$4N$ states / band

$N = 5 \times 10^{22} / \text{cm}^3$

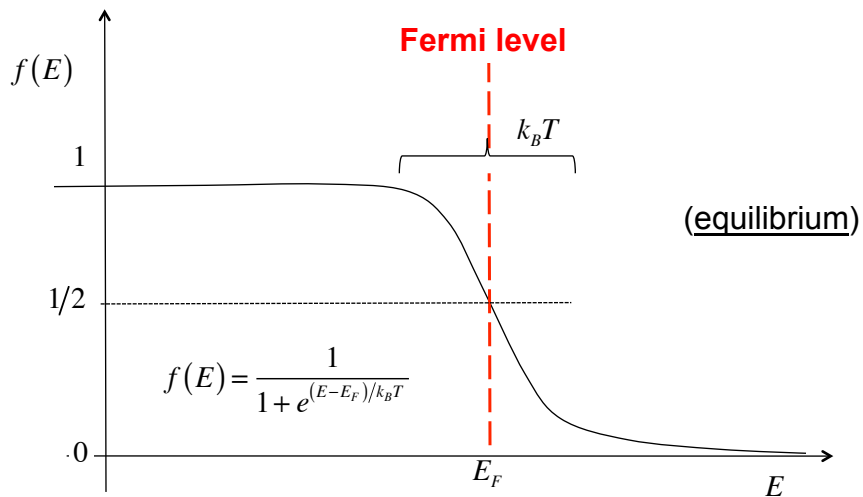
occupying the bands



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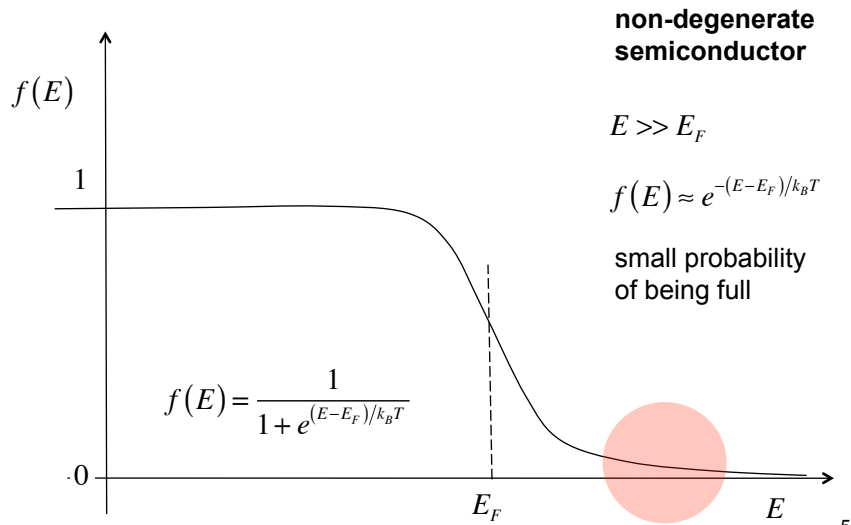
Fermi function



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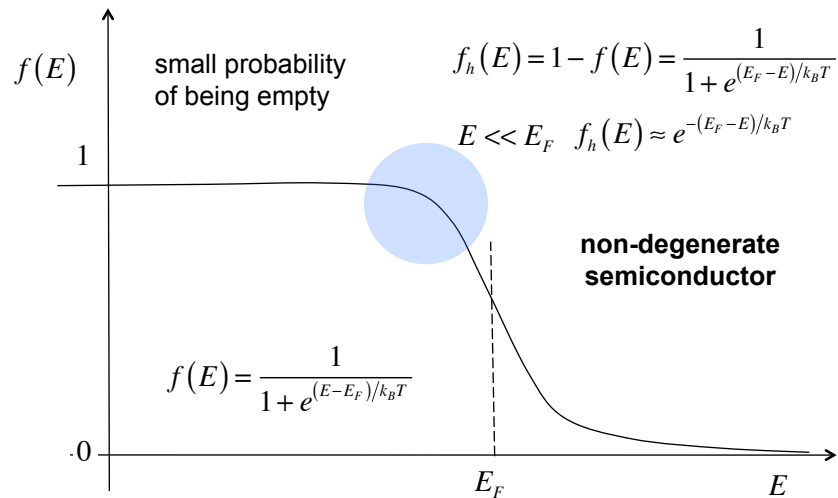
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conduction band



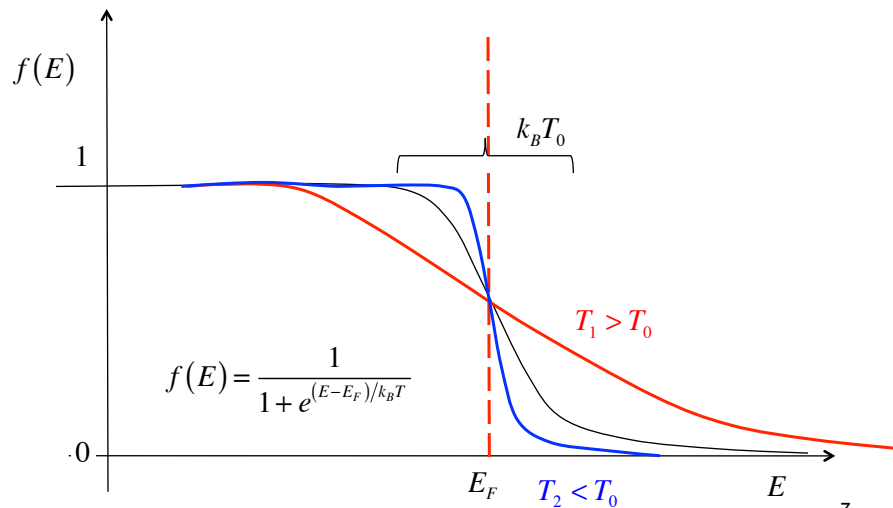
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valence band



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effect of temperature



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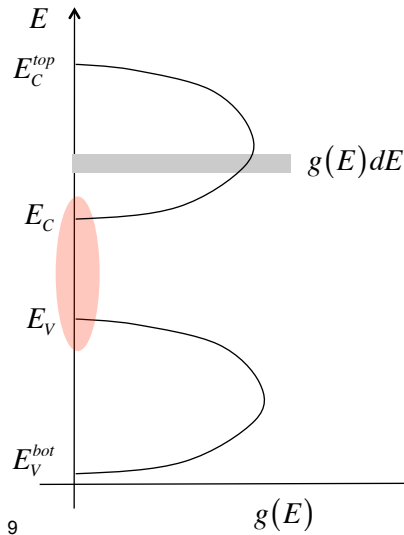
Fermi function

$$f(E) = \frac{1}{1 + e^{(E-E_F)/k_B T}}$$

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density of states

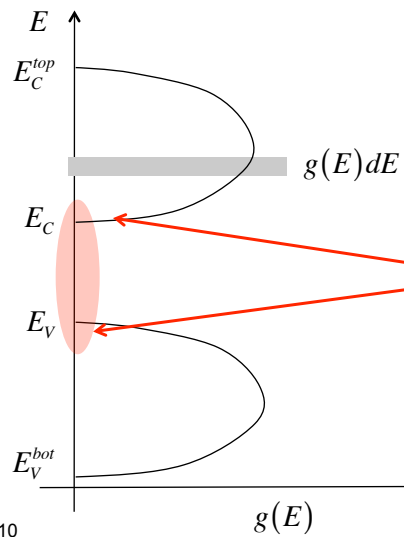


Number of states in an energy range, dE

Units: $\#/J\text{-m}^3$

$$\int_{E_C}^{E_C^{top}} g(E)dE = 4N$$

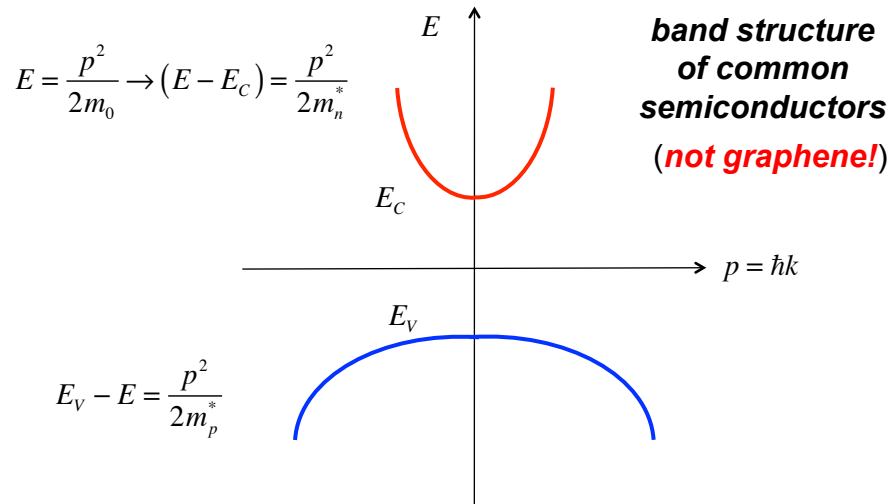
density of states near the band edge



In equilibrium (and near equilibrium), we only need the density of states near the band edge.

Determined by $E(k)$ near the band edge.

$E(k)$ near the band edges



density of states in k-space

1D:

$$N_k dk = 2 \times \left(\frac{L}{2\pi} \right) dk = \frac{L}{\pi} dk$$

2D:

$$N_k d^2k = 2 \times \frac{A}{(2\pi)^2} d^2k = \frac{A}{2\pi^2} d^2k$$

3D:

$$N_k d^3k = 2 \times \left(\frac{\Omega}{8\pi^2} \right) d^3k = \frac{\Omega}{4\pi^3} d^3k$$

Things are simple
in k-space!

density of states in k-space

$$n_S = \frac{1}{A} \sum_{\vec{k}} f_0(E_k) \rightarrow \int_{BZ} f_0(E_k) N_k dk_x dk_y \quad \text{cm}^{-2}$$

$$N_k = 2 \times \left(\frac{A}{4\pi^2} \right) = \frac{A}{2\pi^2}$$

proportional to area

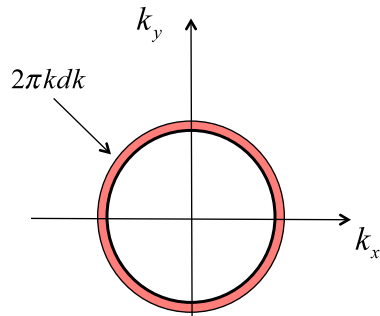
density of states in energy space

$$n_S = \frac{1}{A} \sum_{\vec{k}} f_0(E_k) \rightarrow \int_{E_{BOT}}^{E_{TOP}} f_0(E_k) D_{2D}(E) dE$$

$$\frac{N_k dk}{A} = D_{2D}(E) dE$$

per unit area

2D DOS(E) – parabolic bands



$$E = \frac{\hbar^2 k^2}{2m^*} \quad dE = \frac{\hbar^2 k dk}{m^*}$$

$$N_{2D}(k) dk = \left(\frac{A}{(2\pi)^2} \times 2 \right) dk_x dk_y$$

$$D_{2D}(E) dE = N_{2D}(k) 2\pi k dk / A$$

$$D_{2D}(E) dE = \left(\frac{1}{2\pi^2} \right) 2\pi k dk$$

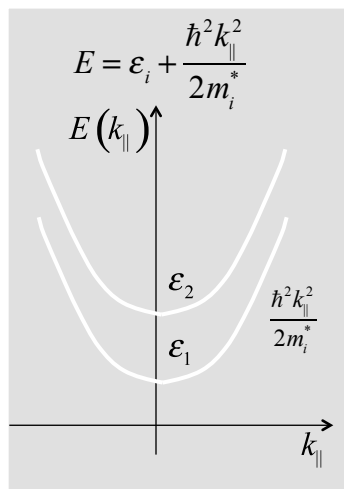
$$D_{2D}(E) dE = \frac{m^*}{\pi \hbar^2} dE$$

$$D_{2D}(E) = \frac{m^*}{\pi \hbar^2}$$

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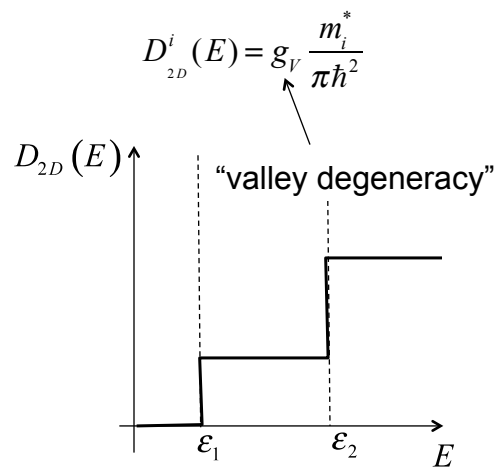
2D DOS(E) – subbands



$$E = \varepsilon_i + \frac{\hbar^2 k_{\parallel}^2}{2m_i^*}$$

$$E(k_{\parallel})$$

$$\frac{\hbar^2 k_{\parallel}^2}{2m_i^*}$$



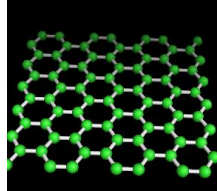
$$D_{2D}^i(E) = g_V \frac{m_i^*}{\pi \hbar^2}$$

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2D DOS(E) – graphene

Graphene is a one-atom-thick planar carbon sheet with a honeycomb lattice.



source: CNTBands 2.0 on nanoHUB.org

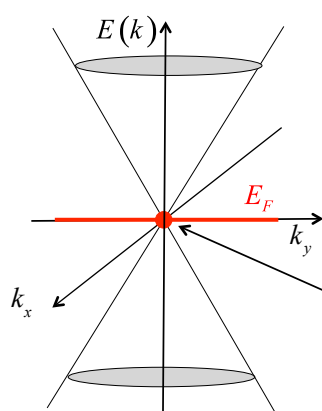
Graphene has an unusual bandstructure that leads to interesting effects and potentially to useful electronic devices.

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graphene: simplified E(k)

We will use a very simple description of the graphene bandstructure, which is a good approximation near the Fermi level.



$$E(k) = \pm \hbar v_F k = \pm \hbar v_F \sqrt{k_x^2 + k_y^2}$$

$$v(k) = \frac{1}{\hbar} \frac{\partial E}{\partial k} = v_F \approx 1 \times 10^8 \text{ cm/s}$$

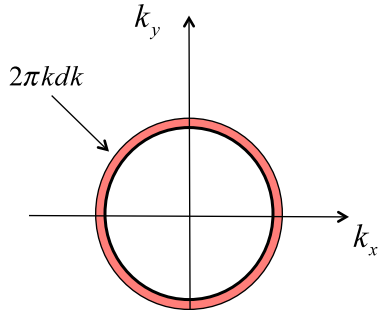
$$g_v = 2 \quad (\text{valley degeneracy})$$

“neutral point” (“Dirac point”)

We will refer to the $E_F > 0$ case, as “n-type graphene” and to the $E_F < 0$ case as “p-type graphene.” 18

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2D DOS(E) – graphene



$$N_{2D}(k)dk = \left(\frac{A}{(2\pi)^2} \times 2 \right) dk_x dk_y$$

$$D_{2D}(E)dE = N_{2D}(k)2\pi kdk / A$$

$$D_{2D}(E)dE = \frac{1}{2\pi^2} 2\pi kdk$$

$$D_{2D}(E)dE = \frac{E}{\pi(\hbar v_F)^2} dE$$

$$E(k) = \hbar v_F k \quad dE = \hbar v_F dk$$

$$k = \frac{E}{\hbar v_F} \quad kdk = \frac{E dE}{(\hbar v_F)^2}$$

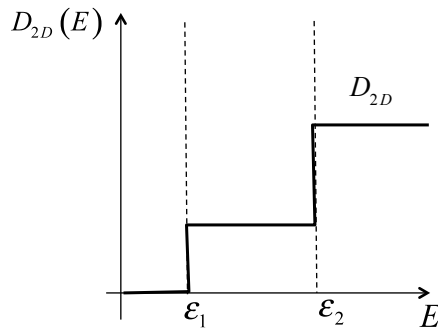
$$D_{2D}(E) = g_V \frac{E}{\pi(\hbar v_F)^2}$$

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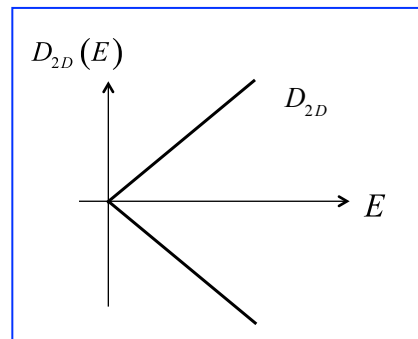
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2D DOS (E)

parabolic



graphene



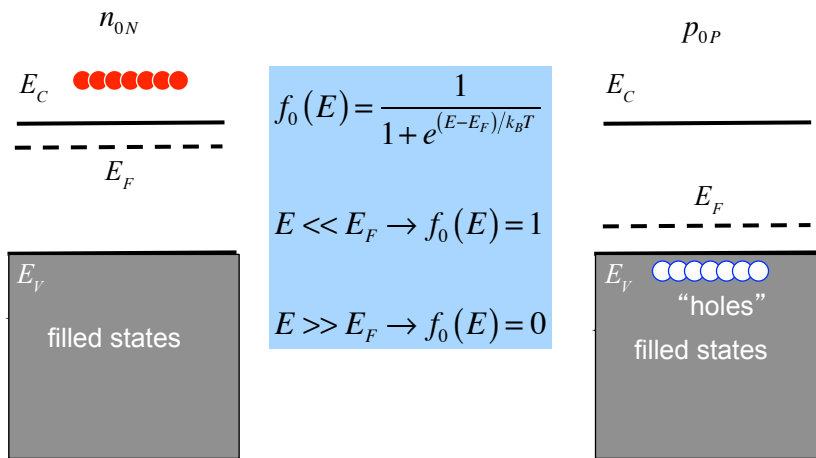
$$D_{2D}^i(E) = g_V \frac{m_i^*}{\pi \hbar^2}$$

$$D_{2D}(E) = g_V \frac{|E|}{\pi \hbar^2 v_F^2} = \frac{2|E|}{\pi \hbar^2 v_F^2}$$

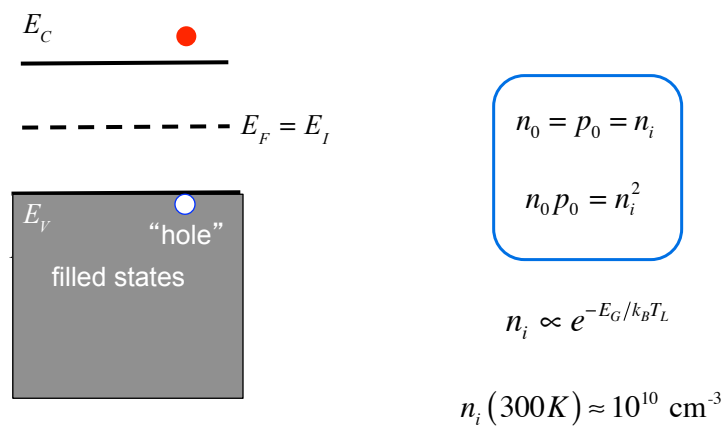
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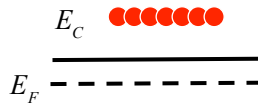
Fermi level (electrochemical potential)



intrinsic semiconductor (3D)

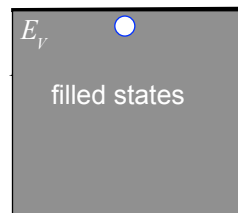


n-type semiconductor (3D)

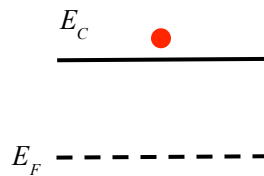


Expect:

$$n_0 = N_C e^{(E_F - E_C)/k_B T}$$

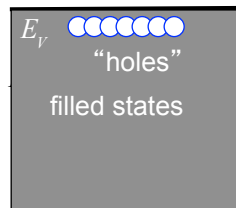


p-type semiconductor (3D)

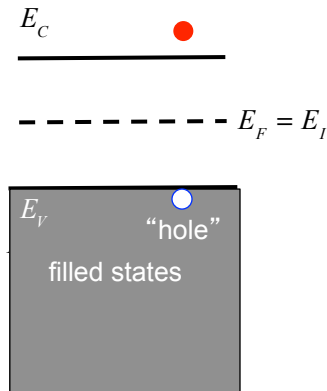


Expect:

$$p_0 = N_V e^{(E_V - E_F)/k_B T}$$



intrinsic semiconductor



$$n_0 = p_0 = n_i$$

$$n_0 p_0 = n_i^2$$

$$n_0 = N_C e^{(E_F - E_C)/k_B T}$$

$$p_0 = N_V e^{(E_V - E_F)/k_B T}$$

$$n_i^2 = N_C N_V e^{-E_G/k_B T}$$

Example: 2D parabolic energy bands

$$n_S = \int_{E_C}^{E_{TOP}} f_0(E) D_{2D}(E) dE \text{ cm}^{-2}$$

$$n_S = \int_{E_C}^{E_{TOP}} \frac{1}{1 + e^{(E - E_F)/k_B T}} \left(g_V \frac{m_n^*}{\pi \hbar^2} \right) dE \text{ cm}^{-2}$$

$$n_S = \left(g_V \frac{m_n^*}{\pi \hbar^2} \right) \int_{E_C}^{\infty} \frac{dE}{1 + e^{(E - E_F)/k_B T}}$$

$$n_S = \left(g_V \frac{m_n^*}{\pi \hbar^2} \right) \int_{E_C}^{\infty} \frac{dE}{1 + e^{(E - E_C + E_C - E_F)/k_B T}}$$

$$f_0(E) = \frac{1}{1 + e^{(E - E_F)/k_B T}}$$

$$D_{2D}(E) = g_V \frac{m_n^*}{\pi \hbar^2}$$

$$\eta \equiv (E - E_C)/k_B T$$

$$\eta_F \equiv (E_F - E_C)/k_B T$$

Example: 2D parabolic energy bands

$$n_S = \left(g_V \frac{m_n^*}{\pi \hbar^2} \right) \int_{E_C}^{\infty} \frac{dE}{1 + e^{(E - E_C + E_C - E_F)/k_B T}}$$

$$\eta \equiv (E - E_C)/k_B T$$

$$dE = k_B T d\eta$$

$$n_S = \left(\frac{g_V m_n^* k_B T}{\pi \hbar^2} \right) \int_0^{\infty} \frac{d\eta}{1 + e^{\eta - \eta_F}}$$

$$\eta_F \equiv (E_F - E_C)/k_B T$$

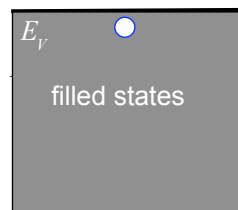
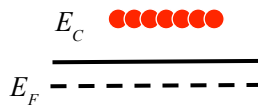
$$\int_0^{\infty} \frac{d\eta}{1 + e^{\eta - \eta_F}} = \ln(1 + e^{\eta_F}) = \mathcal{F}_0(\eta_F)$$

$$n_S = \left(\frac{g_V m_n^* k_B T}{\pi \hbar^2} \right) \mathcal{F}_0(\eta_F) = N_C^{2D} \mathcal{F}_0(\eta_F)$$

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N-type semiconductor (2D parabolic)



$$n_S = N_C^{2D} \mathcal{F}_0(\eta_F)$$

$$N_C^{2D} = \left(\frac{g_V m_n^* k_B T}{\pi \hbar^2} \right)$$

$$\mathcal{F}_0(\eta_F) = \ln(1 + e^{\eta_F})$$

$$\eta_F \ll 0 \quad E_F \ll E_C :$$

$$\mathcal{F}_0(\eta_F) \rightarrow e^{\eta_F}$$

$$n_S = N_C^{2D} e^{\eta_F} \quad \text{non-degenerate semiconductor}$$

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Example: 2D linear energy bands

$$n_S = \int_0^{E_{TOP}} f_0(E) D_{2D}(E) dE \text{ cm}^{-2}$$

$$n_S = \int_{E_C}^{E_{TOP}} \frac{1}{1 + e^{(E-E_F)/k_B T}} \left(\frac{2E}{\pi \hbar^2 v_F^2} \right) dE \text{ cm}^{-2}$$

$$n_S = \left(\frac{2}{\pi \hbar^2 v_F^2} \right) \int_0^{\infty} \frac{E dE}{1 + e^{(E-E_F)/k_B T}}$$

$$n_S = \left(\frac{2}{\pi \hbar^2 v_F^2} \right) \int_0^{\infty} \frac{k_B T \eta (k_B T d\eta)}{1 + e^{\eta - \eta_F}}$$

$$f_0(E) = \frac{1}{1 + e^{(E-E_F)/k_B T}}$$

$$D_{2D}(E) = \frac{2E}{\pi \hbar^2 v_F^2}$$

$$\eta \equiv E/k_B T$$

$$\eta_F \equiv E_F/k_B T$$

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Example: 2D linear energy bands

$$n_S = \frac{2}{\pi} \left(\frac{k_B T}{\hbar v_F} \right)^2 \int_0^{\infty} \frac{\eta d\eta}{1 + e^{\eta - \eta_F}}$$

$$\mathcal{F}_1(\eta_F) \equiv \int_0^{\infty} \frac{\eta d\eta}{1 + e^{\eta - \eta_F}}$$

$$n_S = N_C^{\text{graphene}} \mathcal{F}_1(\eta_F) \text{ cm}^{-2}$$

$$N_C^{\text{graphene}} = \frac{2}{\pi} \left(\frac{k_B T}{\hbar v_F} \right)^2$$

$$\eta \equiv E/k_B T$$

$$\eta_F \equiv E_F/k_B T$$

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Fermi-Dirac integrals

$$\mathcal{F}_j(\eta_F) = \frac{1}{\Gamma(j+1)} \int_0^\infty \frac{\eta^j d\eta}{1 + e^{\eta - \eta_F}}$$

$$\Gamma(n) = (n-1)! \quad (n \text{ integer})$$

$$\Gamma(1/2) = \sqrt{\pi}$$

$$\Gamma(p+1) = p\Gamma(p)$$

$$\mathcal{F}_j(\eta_F) \rightarrow e^{\eta_F} \quad \eta \ll 1$$

$$(E_F - E_C)/k_B T \ll 1$$

$$\frac{d\mathcal{F}_j}{d\eta_F} = \mathcal{F}_{j-1}$$

don't confuse with.... $F_j(\eta) = \int_0^\infty \frac{x^j dx}{1 + e^{x-\eta}}$

For an introduction to Fermi-Dirac integrals, see: "Notes on Fermi-Dirac Integrals," 3rd Ed., by R. Kim and M. Lundstrom) <https://www.nanohub.org/resources/5475>

exercises

$$u = \frac{\langle E - E_C \rangle}{n_S} = \frac{\int_{E_C}^\infty (E - E_C) f_0(E) D_{2D}(E) dE}{\int_{E_C}^\infty f_0(E) D_{2D}(E) dE} \quad \text{J}$$

- 1) $T = 0$ K
- 2) Under non-degenerate conditions
- 3) In general