

Cyclotron emission

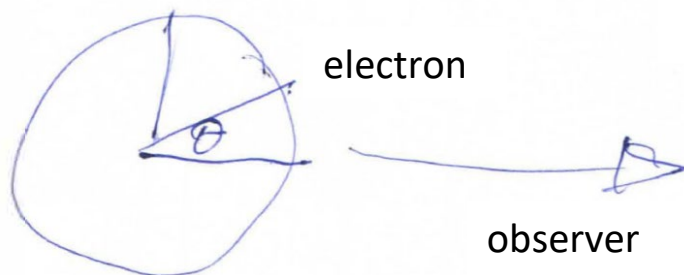
Electron spiraling with the angle α around a magnetic field line moves with acceleration a

$$\vec{F} = -e \vec{v} \times \vec{B}_0 \quad v_{\perp} = v \sin \alpha \quad |\vec{a}| = \left| \frac{d\vec{v}}{dt} \right| = \frac{e v_{\perp} B_0}{m_e} = \omega_c v \sin \alpha$$

Power of emitted radiation is given by the Larmor formula

$$P = \frac{2}{3} \frac{e^2 a^2}{4\pi\epsilon_0 c^3} = \frac{2}{3} \frac{e^2}{c} \frac{e^2 B_0^2}{4\pi\epsilon_0 m_e^2} \frac{v^2 \sin^2 \alpha}{c^2}$$

Observation in the plane $\perp B_0$ (for simplicity)



angular distribution of emission $P_{\Omega} \sim \sin^2 \theta$

time dependence of incident radiation at observer

position $P(t) \sim \sin^2 \theta = \sin^2 \omega_c t$

Other observation angles – harmonics $n\omega_c$ are observed in addition to the frequency ω_c

Emitted power integrated over full solid angle (isotropic electron distribution assumed)

$$\overline{P} = \int_0^\pi P \, 2\pi \sin \alpha \, d\alpha = \frac{2}{3} \frac{e^2}{c} \frac{\omega_c^2}{4\pi\epsilon_0} \left(\frac{v}{c}\right)^2 \int_0^\pi 2\pi \sin^3 \alpha \, d\alpha$$

angular integral = $8\pi/3$

$$\overline{P} = \frac{e^2}{4\pi\epsilon_0} \frac{4\omega_c^2}{3m_e c^3} \frac{8\pi}{3} \frac{m_e v^2}{2} = \frac{e^2}{4\pi\epsilon_0} \frac{4\omega_c^2}{3m_e c^3} \frac{8\pi}{3} E_{kin}$$

Evolution of electron kinetic energy

$$\frac{d E_{kin}}{d t} = -\overline{P} = -\frac{E_{kin}}{t_0} \qquad t_0 = \frac{9\epsilon_0}{8} \frac{m_e^3 c^3}{e^4 B_0^2} \simeq \frac{0.3}{B_0^2} \text{ [s]}$$

Electron loses its energy due to cyclotron emission relatively fast in a very dilute plasma, where reabsorption is negligible.

We shall derive now the absorption coefficient K_c to estimate reabsorption.

Absorption coefficient

The absorption coefficient will be calculated via principle of detailed balancing
In equilibrium – emission = absorption

$$n_e \bar{P} = K_c \bar{I}_{p\nu_c} \Delta\nu$$

Radiation intensity in equilibrium

$$\bar{I}_{p\nu_c} = 4\pi B_{\nu_c} = \frac{8\pi h \nu_c^3 / c^2}{\exp(\frac{h \nu_c}{k_B T_e}) - 1}$$

As ν_c is of order 10^{11} Hz, $h \nu_c \ll k_B T_e$, intensity obeys Rayleigh-Jeans law

$$\bar{I}_{p\nu_c} \simeq \frac{8\pi \nu_c^2 k_B T_e}{c^2} = \frac{2\omega_c^2 k_B T_e}{\pi c^2}$$

Line width of cyclotron emission is due to Doppler broadening

$$\Delta\nu = \nu_c \frac{v}{c} = \frac{\omega_c}{2\pi} \frac{v}{c} \simeq \frac{\omega_c}{2\pi} \sqrt{\frac{2k_B T_e}{m_e c^2}}$$

Absorption coefficient is then

$$K_c = \frac{n_e e^2}{4\pi\epsilon_0} \frac{2\omega_c^2}{3c} \left(\frac{v}{c}\right) \frac{8\pi}{3} \frac{\pi c^2}{2\omega_c^2 k_B T_e} \frac{2\pi c}{\omega_c v} = \frac{n_e e^2}{4\pi\epsilon_0} \frac{16\pi^3}{9B_0} \sqrt{\frac{2m_e}{k_B T_e}}$$

Numerically $K_c \simeq 8.46 \times 10^4 \frac{1}{B_0} \frac{1}{\sqrt{T_e}} \frac{n_e}{10^{19}} \quad [B]=\text{T}, [T_e]=\text{keV}, [n_e]=\text{m}^{-3}, [K_c]=\text{m}^{-1}$

For $n_e = 10^{19} \text{ m}^{-3}$, $B = 1 \text{ T}$, $T_e = 1 \text{ keV}$ $l_f = K_c^{-1} = 1.18 \times 10^{-5} \text{ m}$

and the energy confinement time is estimated $\tau_{\text{energy}} \cong t_0 R / l_f$,

for $R = 0.1 \text{ m}$ $\tau_{\text{energy}} \cong 0.3 \times 0.1 / 1.18 \times 10^{-5} = 2540 \text{ s}$

Thus, energy losses due to cyclotron emission are negligible in magnetic confinement