

Derivation of thermal energy conservation from kinetics

$$\begin{aligned} \langle \vec{v} \rangle &= \vec{v}_S \equiv \vec{u} & \vec{V} &= \vec{v} - \vec{u} \\ f_S &\equiv f & m_S &\equiv m \\ n_S &\equiv n & T_S &\equiv T \end{aligned}$$

Kinetic equation

$$\frac{\partial f}{\partial t} + \vec{v} \frac{\partial f}{\partial \vec{r}} + \frac{\vec{F}}{m} \frac{\partial f}{\partial \vec{v}} = \left(\frac{\partial f}{\partial t} \right)_c$$

We want to derive an equation for thermal energy

$$\int \frac{m}{2} v^2 d^3V \quad d^3V = d^3v$$

1st term

$$\begin{aligned} \int \frac{m}{2} (\vec{v} - \vec{u})^2 \frac{\partial f}{\partial t} d^3v &= \frac{\partial}{\partial t} \int \frac{m}{2} v^2 f d^3v - \\ - \frac{m}{2} \frac{\partial \vec{u}}{\partial t} \int -2(\vec{v} - \vec{u}) f d^3v &= \frac{\partial}{\partial t} \left(\frac{3}{2} n k_B T \right) \end{aligned}$$

2nd term

$$\begin{aligned} \frac{m}{2} \int v_i \frac{\partial f}{\partial n_i} (v_j - u_j)^2 d^3v &= \frac{\partial}{\partial n_i} \int \frac{m v_i}{2} f (v_j - u_j)^2 d^3v + \\ + \frac{m}{2} 2 \frac{\partial u_j}{\partial n_i} \int d^3v v_i (v_j - u_j)^2 f &= \\ = \frac{\partial}{\partial n_i} \int \frac{m}{2} (v_i - u_i) (\vec{v} - \vec{u})^2 f d^3v + \frac{\partial}{\partial n_i} \int \frac{m u_i}{2} (\vec{v} - \vec{u})^2 f d^3v + \\ + \frac{\partial u_j}{\partial n_i} \int (v_i - u_i) (v_j - u_j) m f d^3v + \frac{\partial u_j}{\partial n_i} u_i \int m (v_j - u_j) f d^3v &= \\ = \nabla \vec{q} + \nabla \left(\vec{u} \frac{3}{2} n k_B T \right) + P_{ij} \frac{\partial u_j}{\partial n_i} + 0 \end{aligned}$$

3rd term

$$\begin{aligned}
 & \frac{m}{2} \int \frac{F_i}{m} \frac{\partial f}{\partial v_i} (\vec{v} - \vec{u})^2 d^3v = \int \frac{\partial}{\partial v_i} \left[\frac{F_i}{2} f (\vec{v} - \vec{u})^2 \right] d^3v - \\
 & - \frac{q}{2} \int \frac{\partial}{\partial v_i} \left[\underbrace{E_i + \epsilon_{ijk} v_j B_k}_0 (\vec{v} - \vec{u})^2 \right] d^3v - \underbrace{q E_i \int (v_i - u_i) f d^3v}_0 - \\
 & - \frac{q}{2} \epsilon_{ijk} B_k \int v_j \frac{\partial}{\partial v_i} (\vec{v} - \vec{u})^2 f d^3v = \\
 & = - \frac{q}{m} \epsilon_{ijk} B_k \int m (v_j - u_j) (v_i - u_i) f d^3v - \frac{q}{m} \epsilon_{ijk} B_k \underbrace{\int u_j (v_i - u_i) f d^3v}_0 = \\
 & = - \frac{q}{m} \epsilon_{ijk} B_k P_{ij} = 0
 \end{aligned}$$

$P_{ij} = P_{ji} \quad \epsilon_{ijk} = -\epsilon_{jik}$

Volume force does not change the internal energy

Collisional term

$$\frac{1}{2} \int m v^2 \left(\frac{\partial f}{\partial t} \right)_c d^3v = \sum_t \int \frac{m_s v_s^2}{2} \left(\frac{\partial f}{\partial t} \right)_c^{(st)} d^3v_s \approx \sum_t \alpha_{ts} (T_t - T_s)$$

ss collisions – energy conservation → 0 st collisions – temperature relaxation

$$\begin{aligned}
 & \frac{\partial}{\partial t} \left(\frac{3}{2} n_s k_B T_s \right) + \text{div} \left(\vec{q}_s + \vec{v}_s \frac{3}{2} n_s k_B T_s \right) + P_{ij}^s \frac{\partial v_{si}}{\partial v_j} = \\
 & = \sum_t \alpha_{ts} (T_t - T_s)
 \end{aligned}$$

$$\vec{q}_s = -\kappa_s \nabla T_s$$

for small temperature gradients